

Simulating the EOR with self-consistent growth of galaxies

Master's thesis presentation

ETH Zürich, University of Zürich

Simulating the Epoch of Reionization

BEoRN

Halo growth

Adaptations

Results

Conclusion

End

References

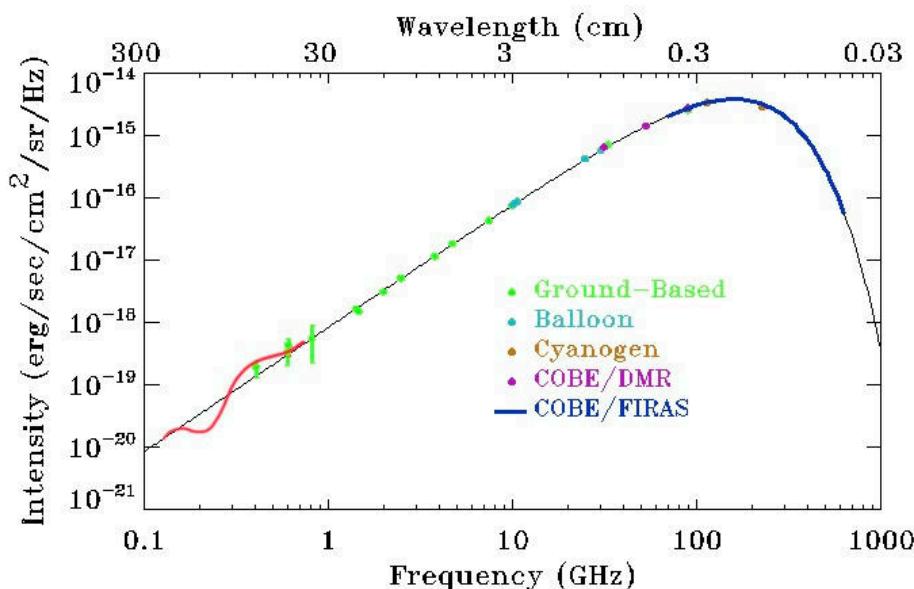
Simulating the Epoch of Reionization

- The 21-cm signal
- Expressing the 21-cm signal [1], [2]
- The current state of simulations

The 21-cm signal

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The *brightness temperature* describes the intensity of the 21-cm line

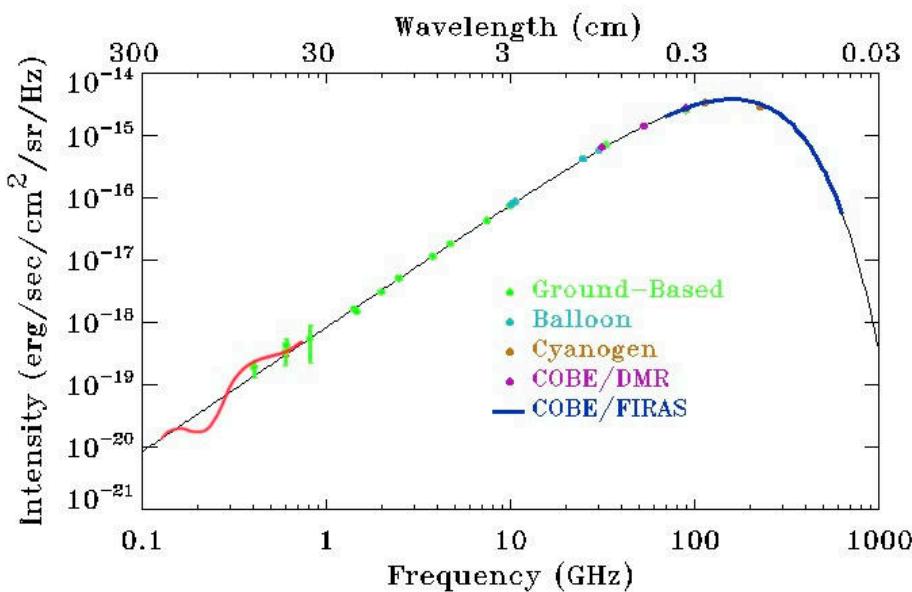


from [3]

The 21-cm signal

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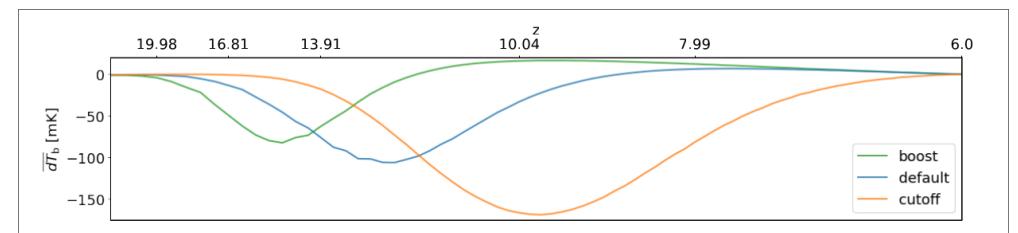
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remove contribution from the BB spectrum

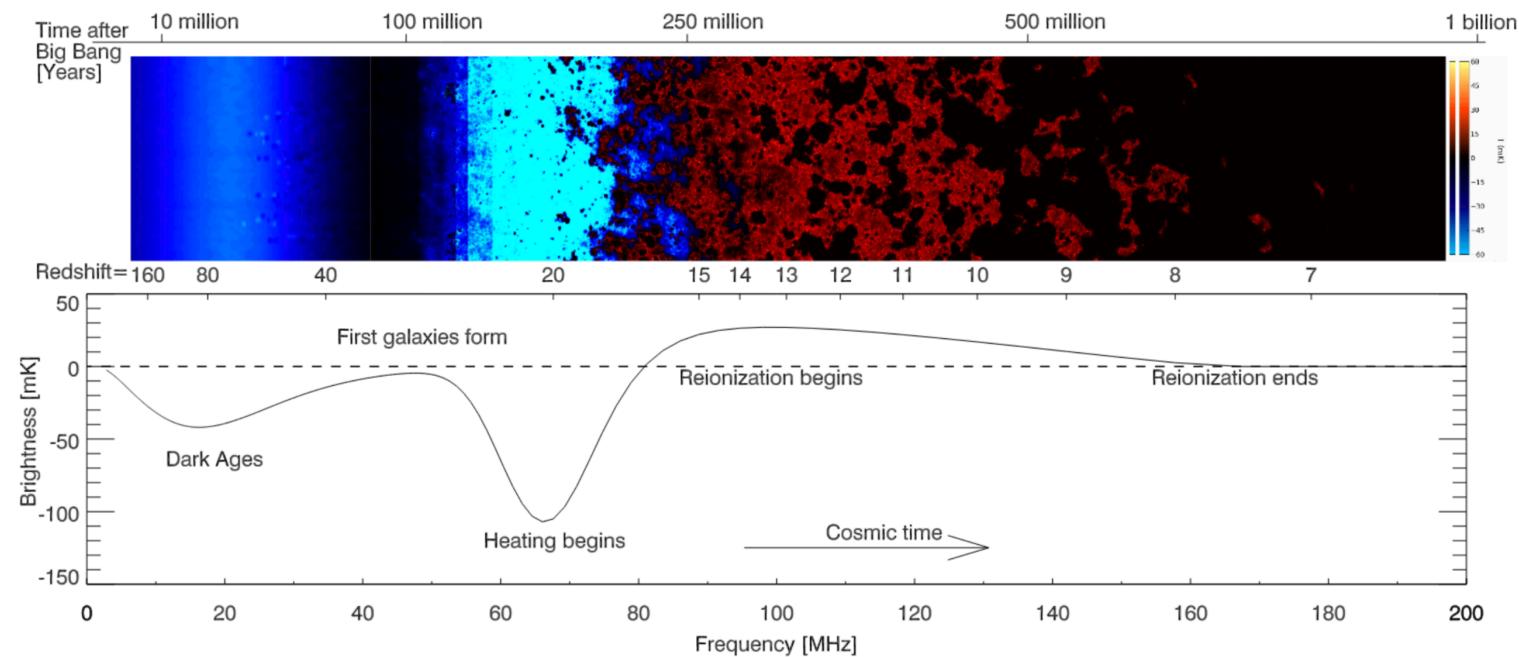
differential brightness temperature

⇒ the actual reionization signal



Expression the 21-cm signal [1], [2]

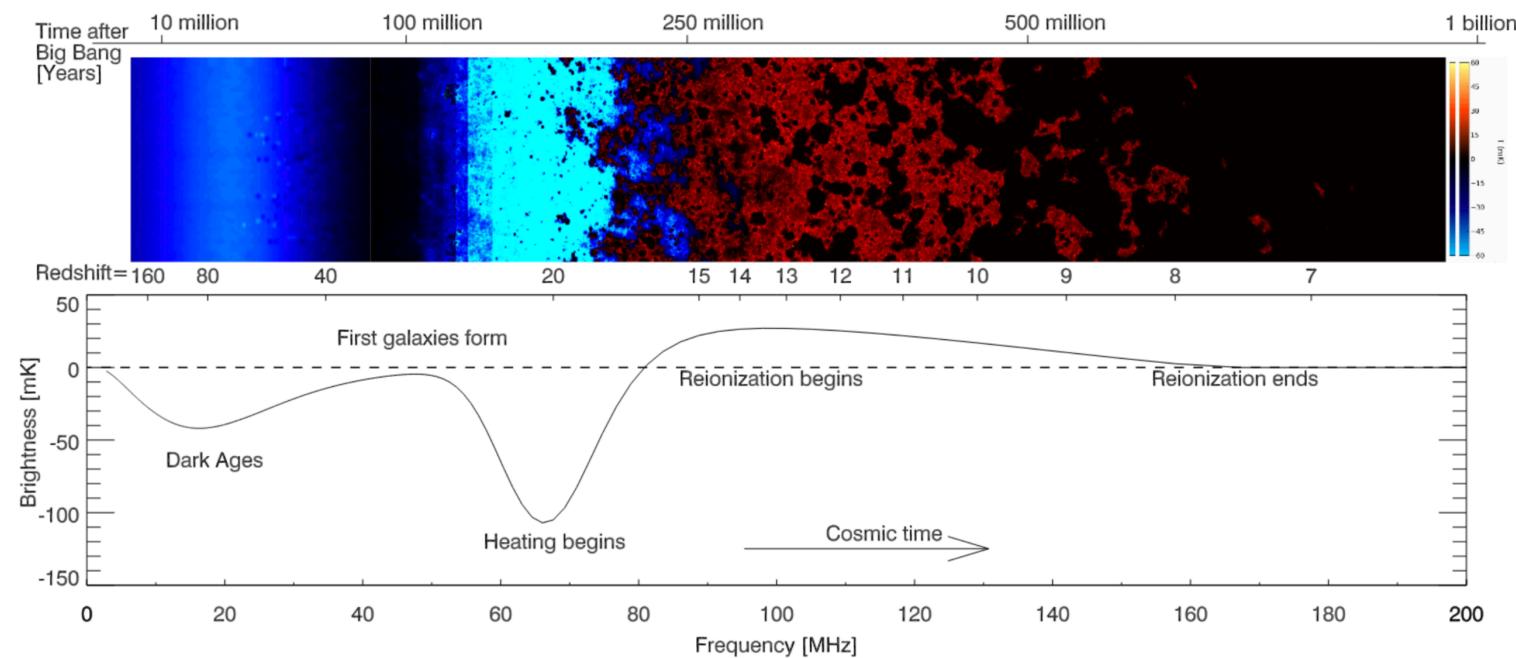
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Simulating the EOR with self-consistent growth of galaxies

Expression the 21-cm signal [1], [2]

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$$dT_b(\mathbf{x}, z) \simeq T_0(z) \cdot x_{\text{HI}}(\mathbf{x}, z) \cdot (1 + \delta_b(\mathbf{x}, z)) \cdot \frac{x_\alpha(\mathbf{x}, z)}{1 + x_\alpha(\mathbf{x}, z)} \cdot \left(\frac{1 - T_{\text{CMB}}(z)}{T_{\text{gas}}(\mathbf{x}, z)} \right)$$

Traditional approaches

- need to cover large dynamic range
- hydrodynamics & radiative transfer
- hard to scale
- ⇒ no reproducibility

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Semi-numerical approaches

- such as BEoRN [4], 21cmFAST [5]
- approximative treatment
- prediction of global signals
- scalable + efficient
- ⇒ reproducible and flexible

Simulating the
Epoch of Reionization

BEoRN

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BEoRN

- The halo model of reionization
- Revisiting the 21cm signal
- The “painting” procedure
- Postprocessing
- Maps
- Signal

Following [6], [7], the halo model describes (derivation):

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$$\rho_\alpha(r \mid M, z) = \frac{(1+z)^2}{4\pi r^2} \cdot \sum_{n=2}^{n_m} f_n \cdot \varepsilon_\alpha(\nu') \cdot f_\star \cdot \dot{M}(z' \mid M, z)$$

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$$\frac{3}{2} \cdot \frac{d\rho_h(r \mid M, z)}{dz} = \frac{3\rho_h(r \mid M, z)}{1+z} - \frac{\rho_{\text{xray}}(r \mid M, z)}{k_B(1+z)H(z)}$$

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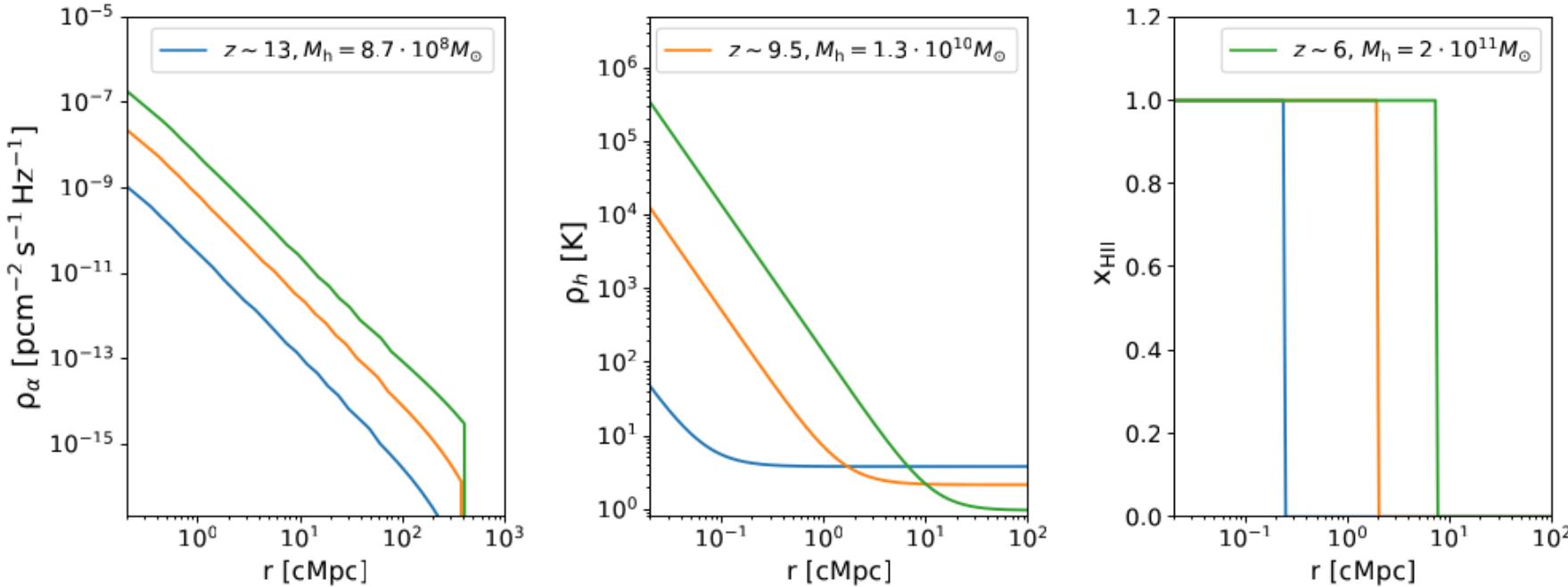
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$$\frac{dV_b}{dt} = \frac{\dot{N}_{\text{ion}(t)}}{\bar{n}_H^0} - \alpha_B \cdot \frac{C}{a^3} \cdot \bar{n}_H^0 \cdot V_b$$

The halo model of reionization

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Visually:



(from [4])

$$dT_b(\mathbf{x}, z) \simeq T_0(z) \cdot x_{\text{HI}}(\mathbf{x}, z) \cdot (1 + \delta_b(\mathbf{x}, z)) \cdot \frac{x_\alpha(\mathbf{x}, z)}{1 + x_\alpha(\mathbf{x}, z)} \cdot \left(\frac{1 - T_{\text{CMB}}(z)}{T_{\text{gas}}(\mathbf{x}, z)} \right)$$

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from x_{HII}

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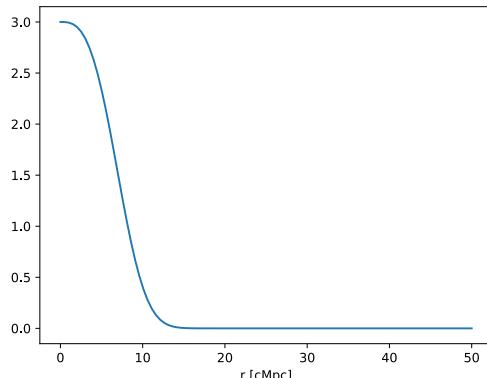
from x_{HII} from ρ_α

$$dT_b(\mathbf{x}, z) \simeq T_0(z) \cdot x_{\text{HI}}(\mathbf{x}, z) \cdot (1 + \delta_b(\mathbf{x}, z)) \cdot \frac{x_\alpha(\mathbf{x}, z)}{1 + x_\alpha(\mathbf{x}, z)} \cdot \left(\frac{1 - T_{\text{CMB}}(z)}{T_{\text{gas}}(\mathbf{x}, z)} \right)$$

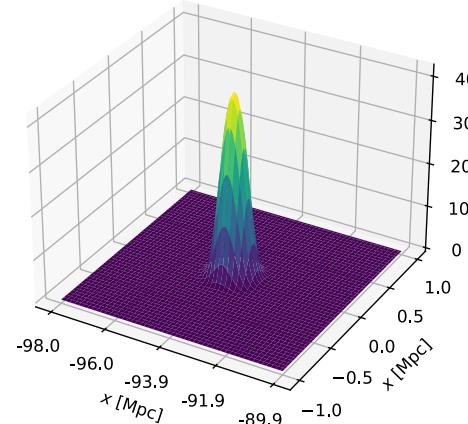
The diagram illustrates the components of the 21cm signal equation. It shows three arrows pointing to specific terms: one arrow points from x_{HII} to the x_{HI} term; another arrow points from ρ_α to the x_α term; and a third arrow points from ρ_h to the T_{gas} term.

The “painting” procedure

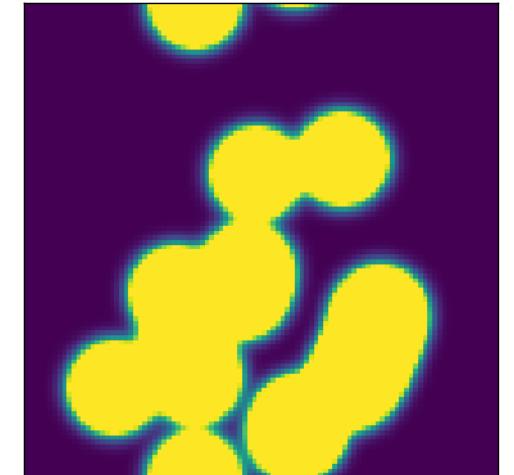
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1-d profile



3-d kernel
(localized)

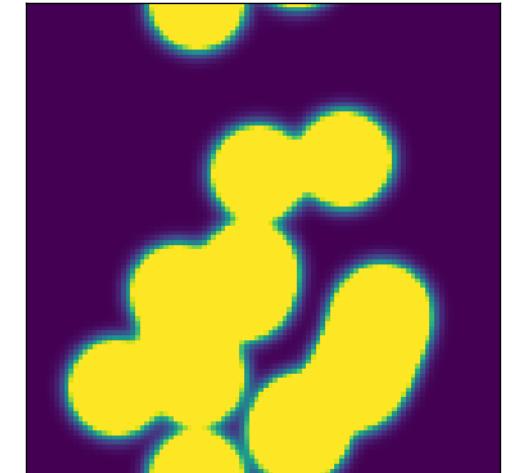
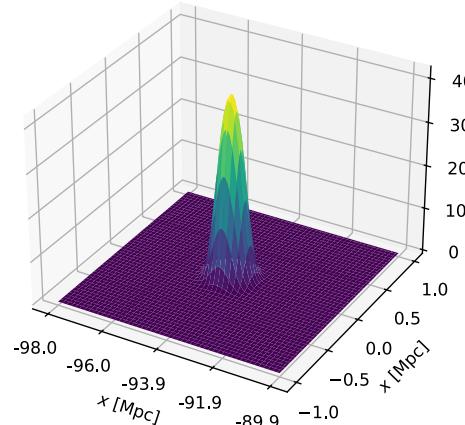
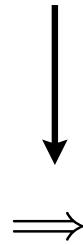
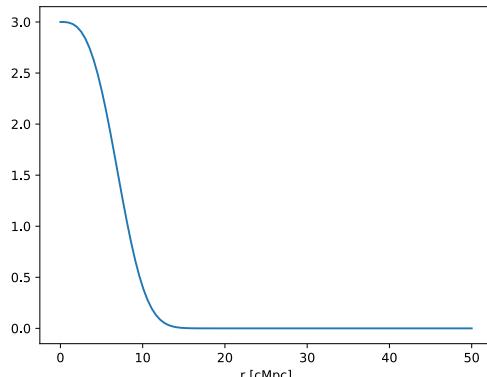


3-d contribution on a
grid

The “painting” procedure

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spherical symmetry



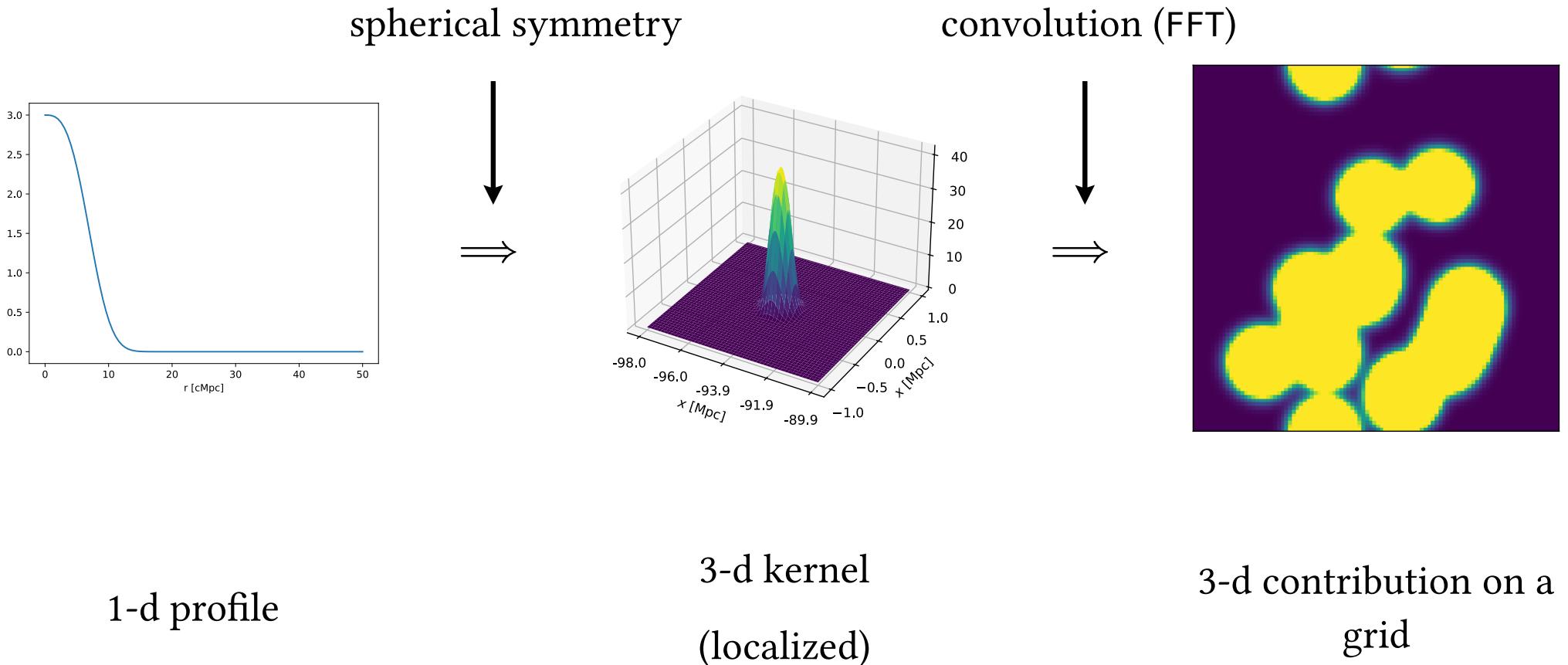
1-d profile

3-d kernel
(localized)

3-d contribution on a
grid

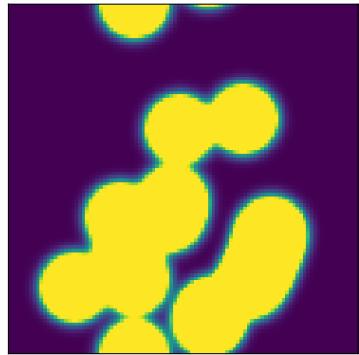
The “painting” procedure

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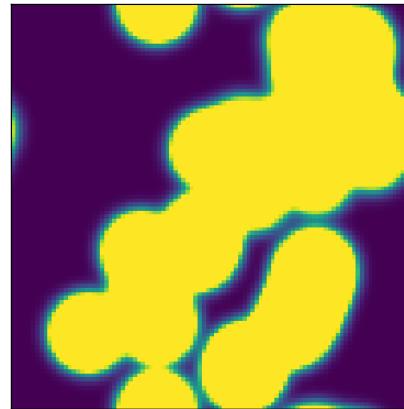
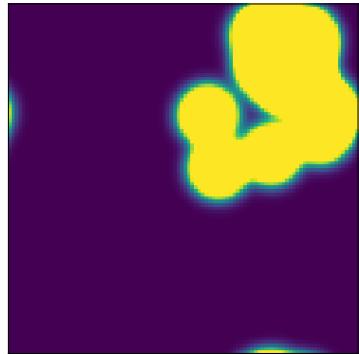


The “painting” procedure

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Multiple contributions \Rightarrow



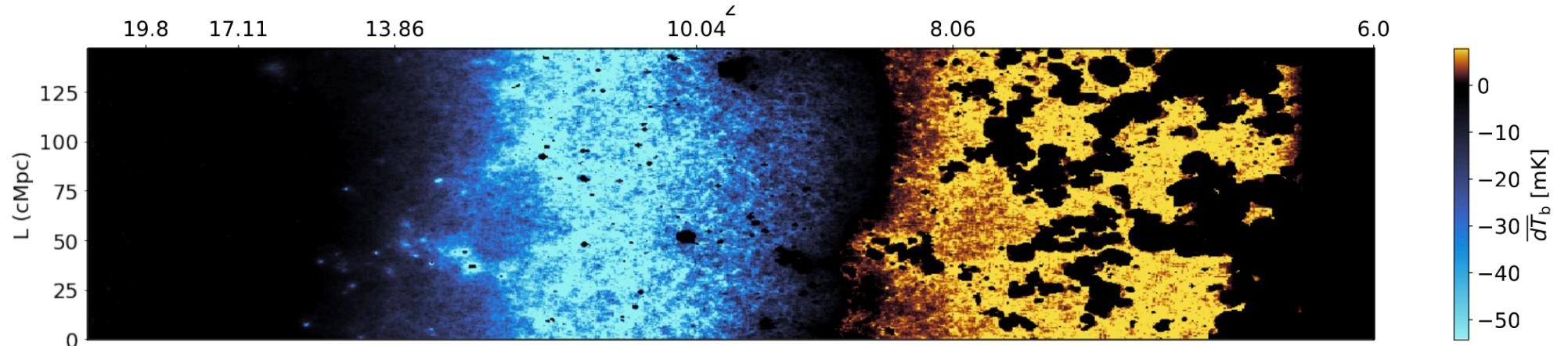
\Rightarrow *Postprocessing*

(overlaps, normalization, ...)

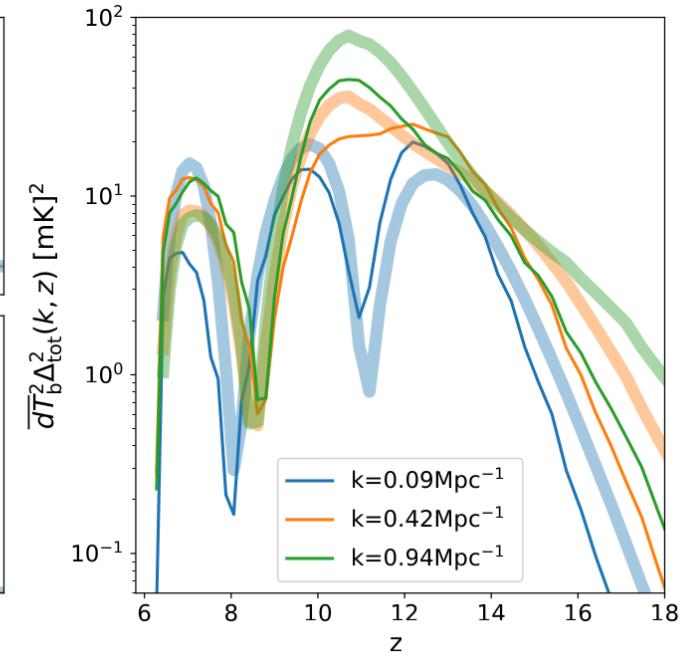
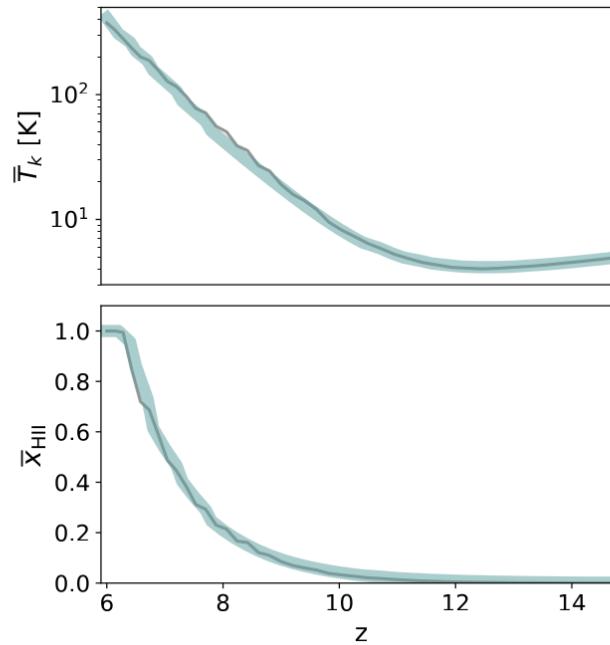
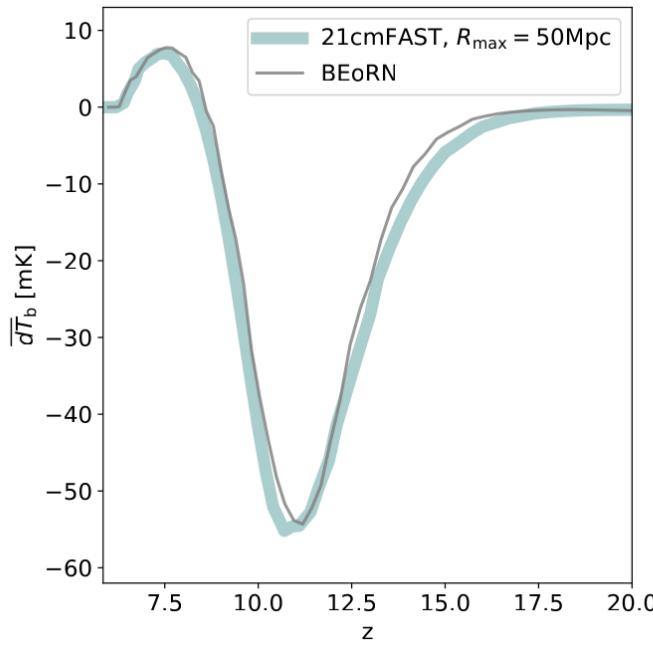
- ionization overlaps
- corrections due to RSD
- computation of derived quantities
- summary statistics

Through the redshifting of photons, the brightness temperature across redshift slices will be measured in a frequency band

⇒ representation as a lightcone



from [4]



from [4]

Simulating the
Epoch of Reionization
BEoRN

Halo growth

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Halo growth

- Motivation
- Effect on the flux profiles
- Inferring growth from THESAN data

Crucial dependence on the **star formation rate**

- assumed to be directly linked to halo growth rate \dot{M} :

$$\dot{M}_\star = f_\star(M_h) \cdot \dot{M}_h$$

- growth according to the exponential model:

$$M_h(z) = M_h(z_0) \cdot \exp[-\alpha(z - z_0)]$$

with $\alpha = \frac{\dot{M}_h}{M_h}$ the *specific growth rate*

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→ **inconsistent** with the N-body output

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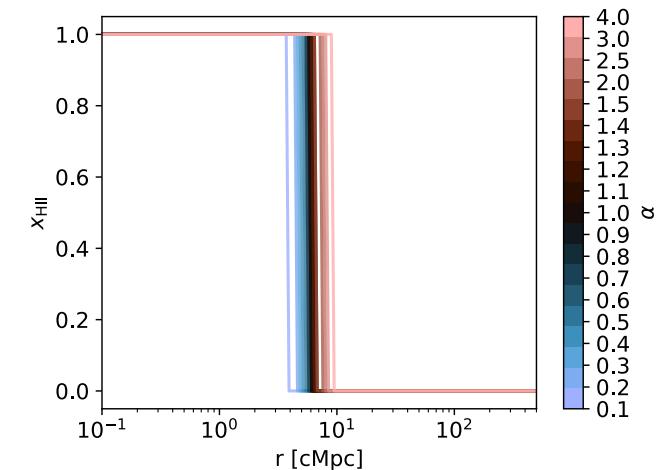
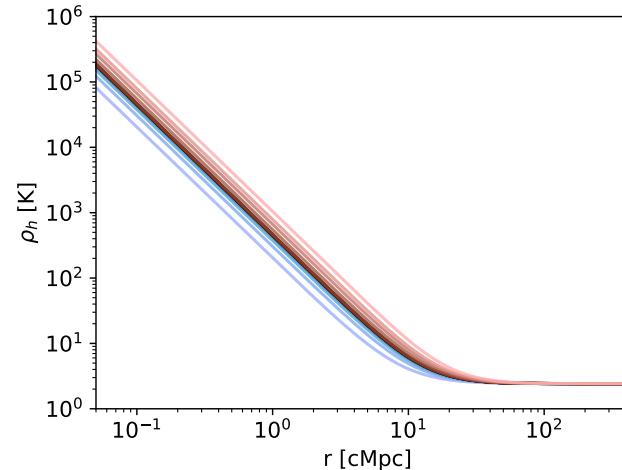
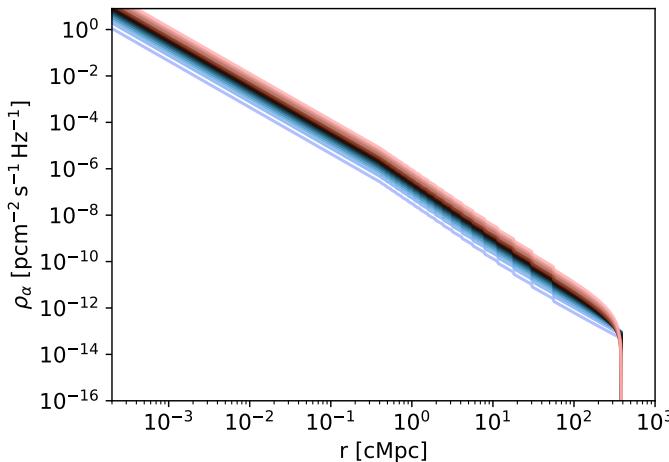
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→ inaccurate when applied to all halos
→ **inconsistent** with the N-body output
→ how to implement **consistent** growth?

Effect on the flux profiles

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$$M_h = 6.08 \cdot 10^{11} M_\odot \text{ (fixed)}$$

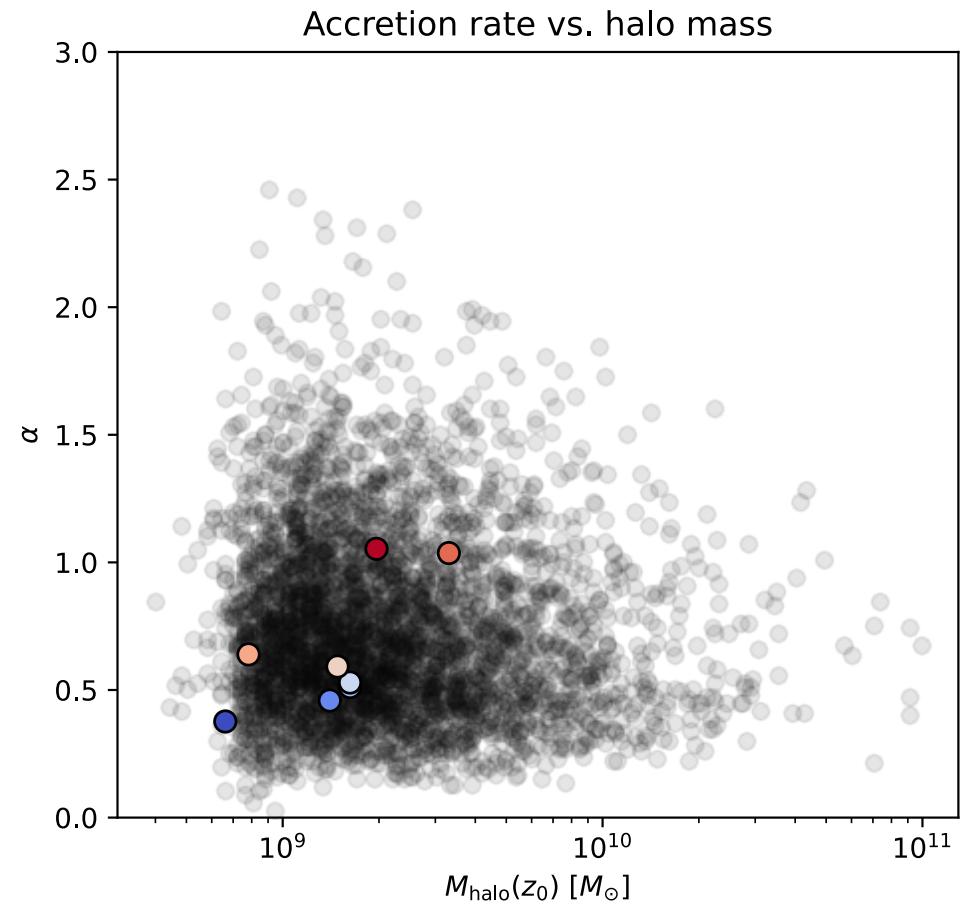
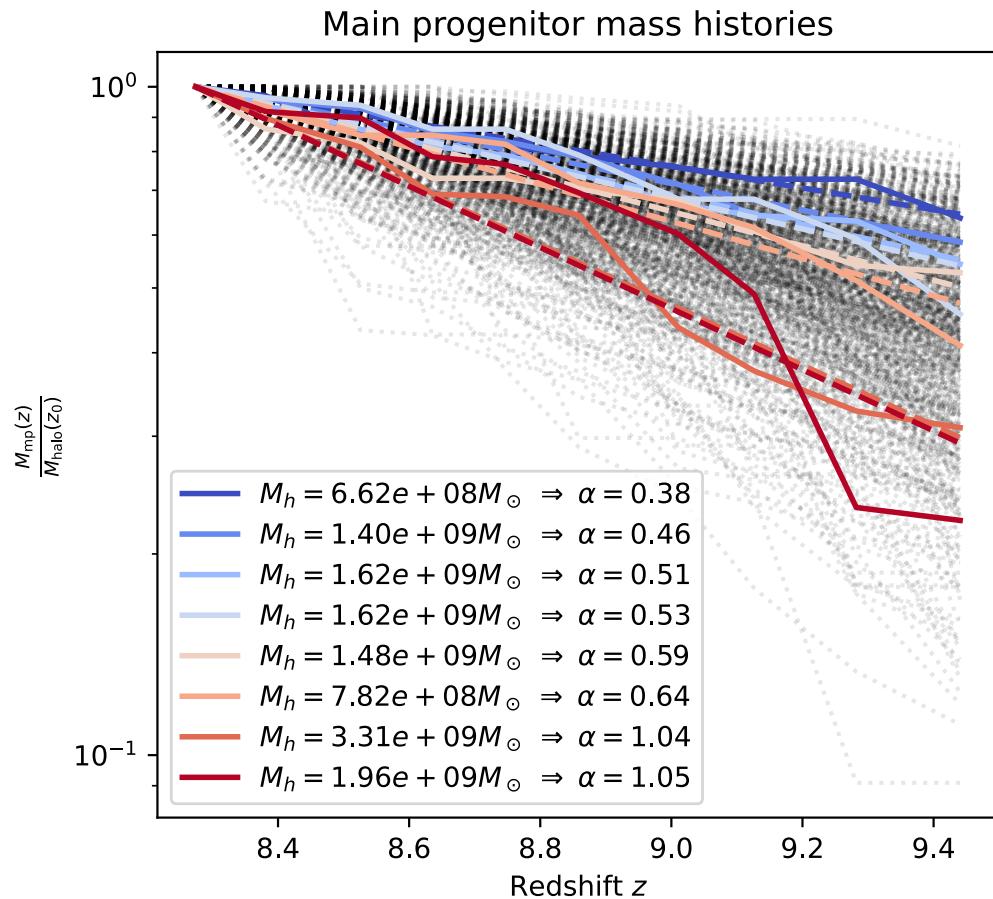
\implies correction up to $\times 5$

- already includes precomputed merger trees [8]
- follow main progenitor branch back in time
- fit the exponential model to main progenitor branch
- use **individual growth** to select profile
- **self-consistent** treatment of halo growth leveraging the snapshots

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Inferring growth from THESAN data

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Adaptations

- Central changes
- Simplified usage

- profile generation taking into account halo growth rate
- reading merger trees + inferring growth rates
- parallel painting across multiple halo bins
- performance and ease of use

- profile generation taking into account halo growth rate
 - reading merger trees + inferring growth rates
 - parallel painting across multiple halo bins
 - performance and ease of use
- Validated 

Simplified usage

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```
#!/usr/bin/env python3

from pathlib import Path
import beorn

current_directory = Path(".")
## Setup the parameters
parameters_file = current_directory / "parameters.yaml"
parameters =
beorn.structs.Parameters.from_yaml(parameters_file)
# sample format:
# parameters.solver.redshifts = [6, 20]
# parameters.simulation.file_root = ... / "Thesan-Dark-1"

## Handling of the io
# this will interface with the input simulation
loader = beorn.load_input_data.ThesanLoader(
    parameters,
    is_high_res = True
)

cache_handler = beorn.io.Handler(current_directory /
"cache")
```

```
output_handler = beorn.io.Handler(current_directory /
"output")
# handlers can also manage logs for us:
# output_handler.save_logs(parameters)

## Computation of the radiation profiles
solver =
beorn.radiation_profiles.ProfileSolver(parameters)
profiles = solver.solve()

## Full 3D painting of the radiation profiles over the
specified redshifts
painter = beorn.painting.Painter(
    parameters,
    loader = loader,
    cache_handler = cache_handler,
    output_handler = output_handler
)

grid = painter.paint_full(profiles)
# Done!
```

Simulating the
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Results

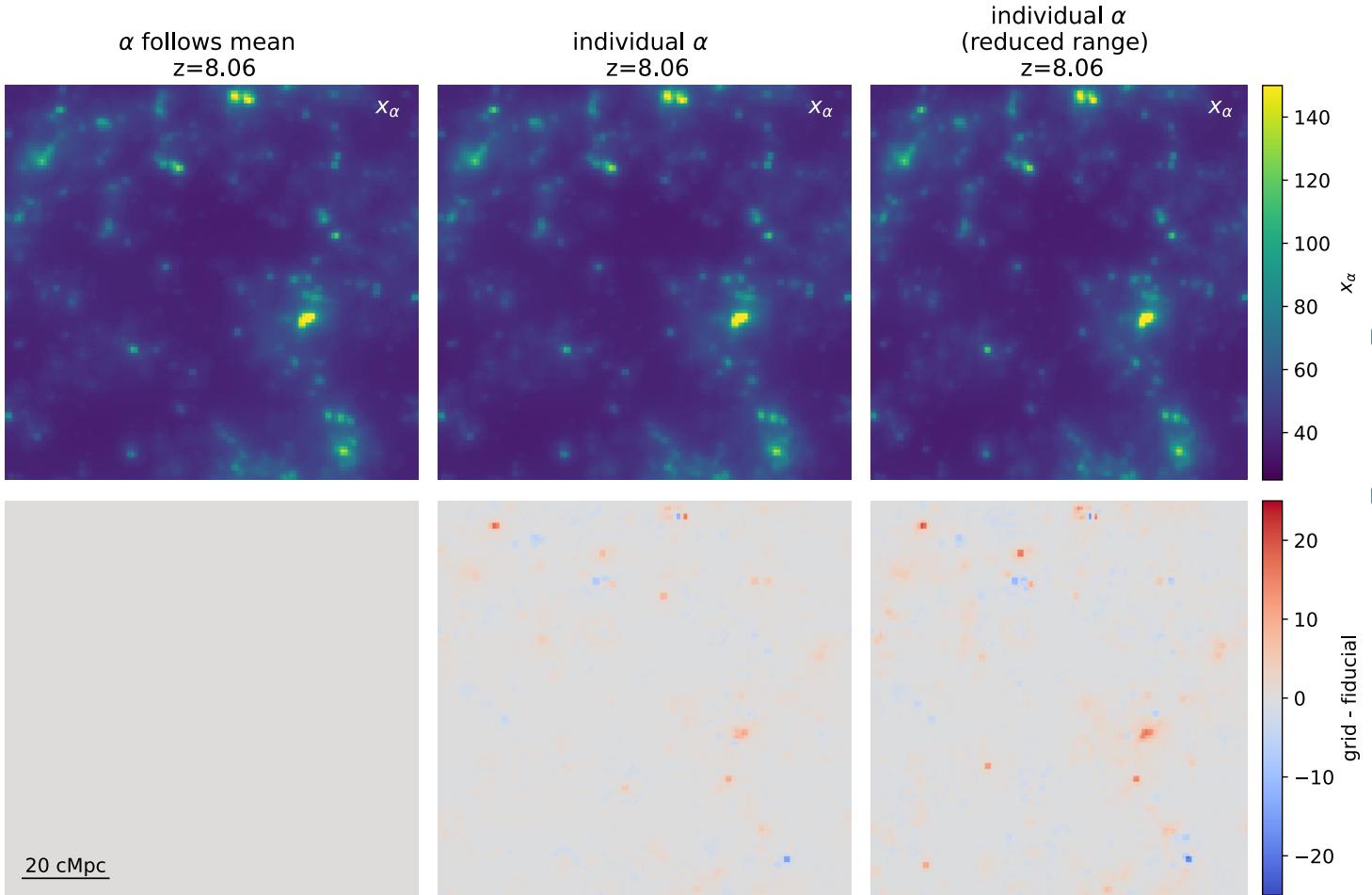
Conclusion
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Results

- Map outputs
- Signals

Map outputs

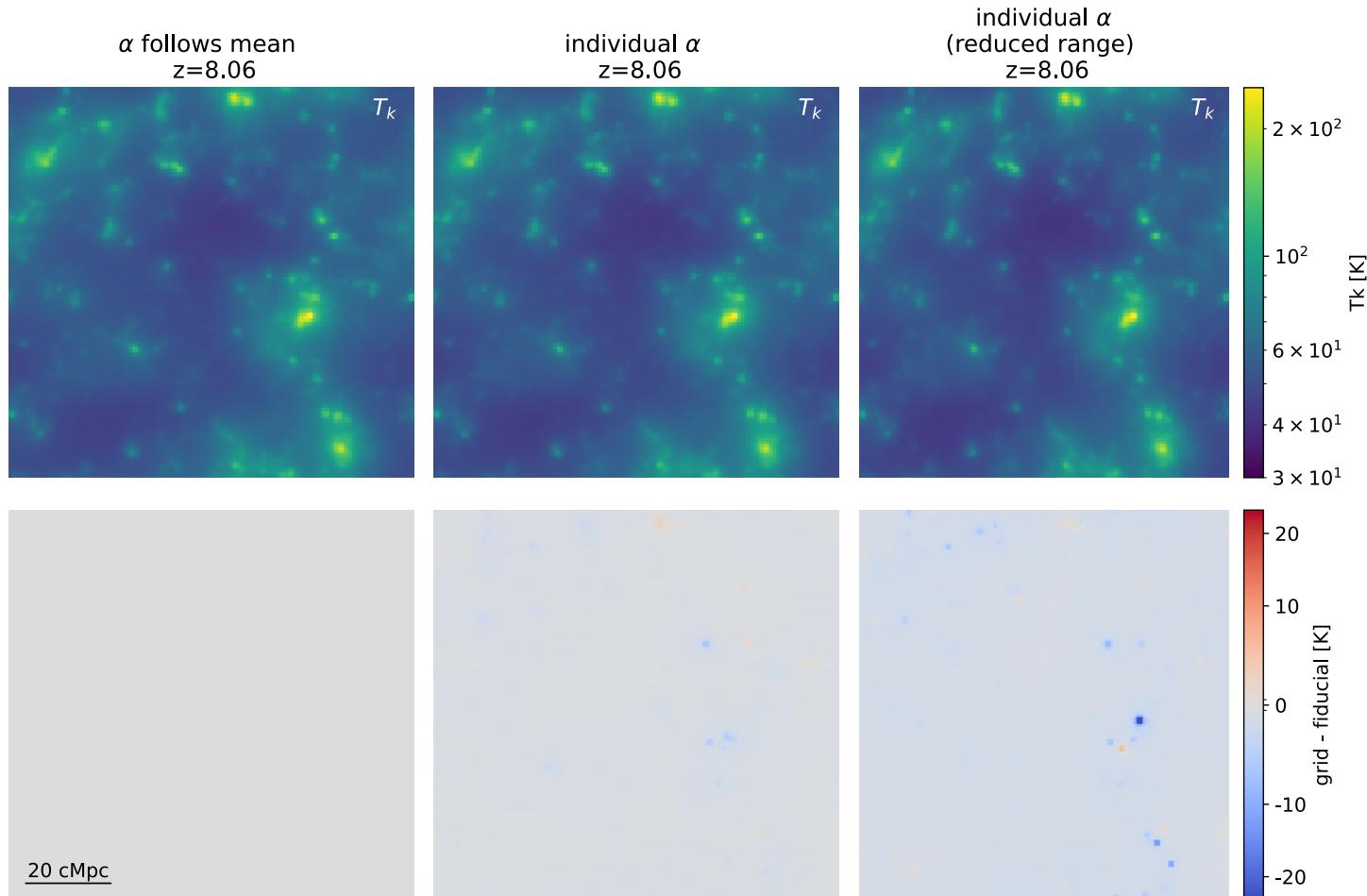
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- strong variations close to sources
- nearly no effect in voids

Map outputs

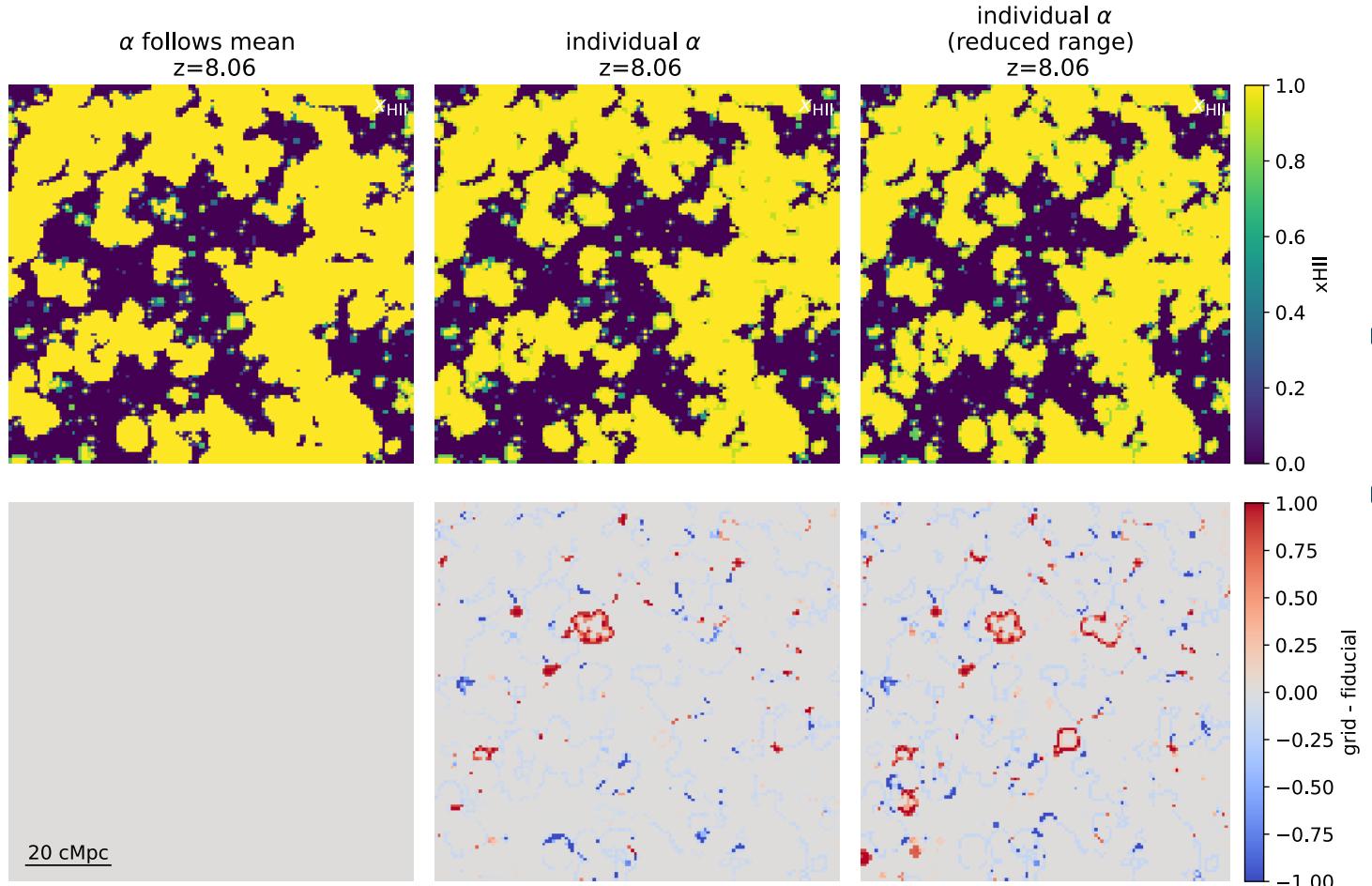
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- delayed heating \Leftrightarrow colder halos
- highest accreting halos catch up

Map outputs

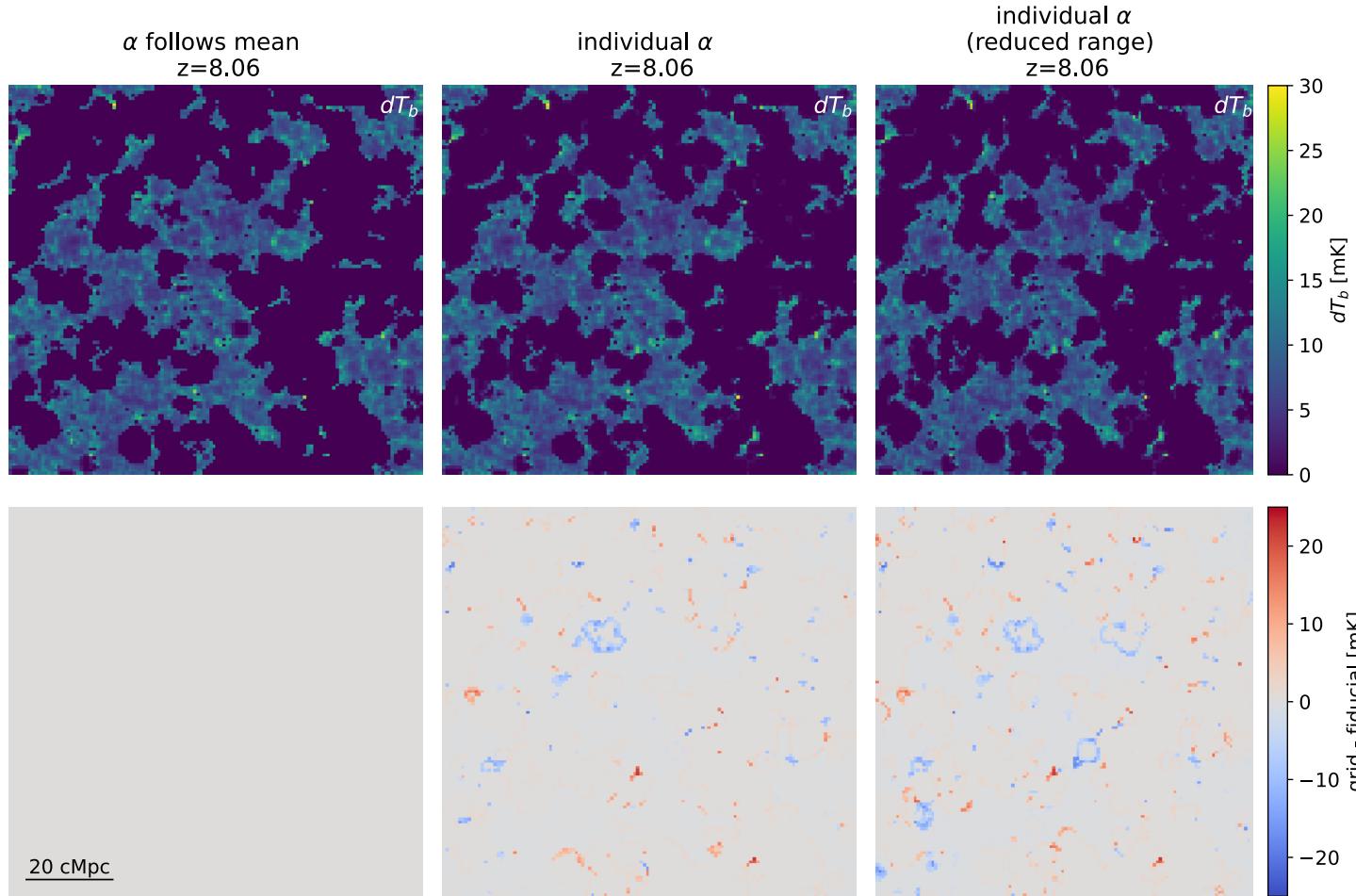
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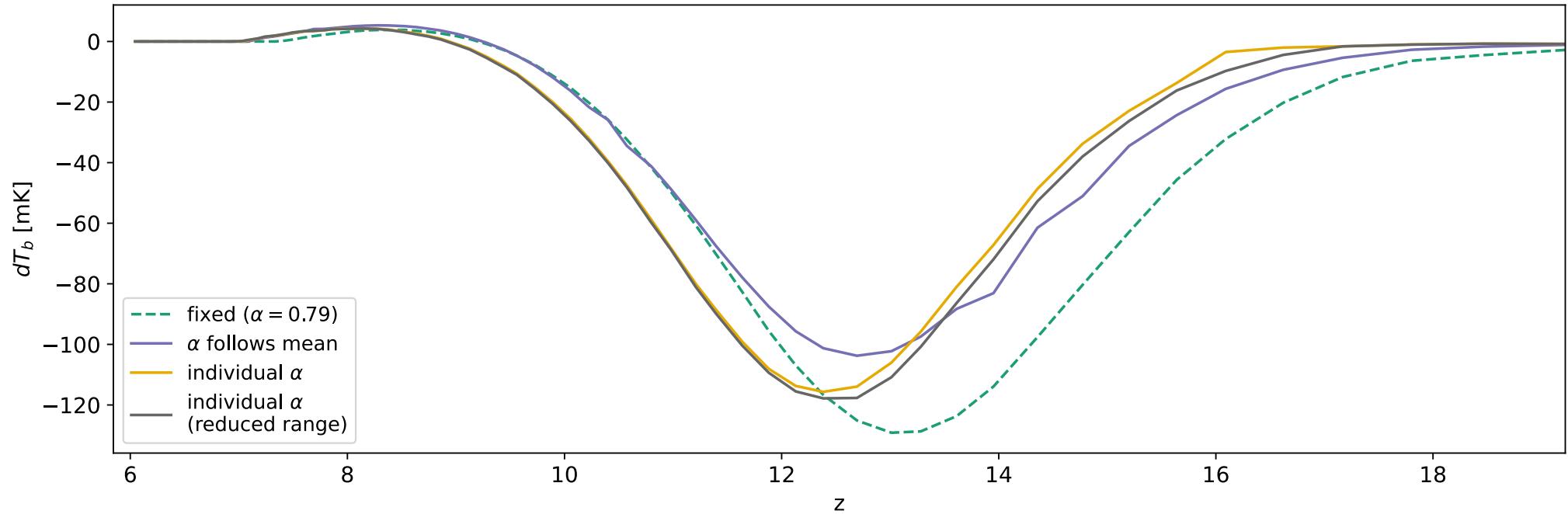
- high contrast due to sharp cutoffs
- clearly increased dynamic range

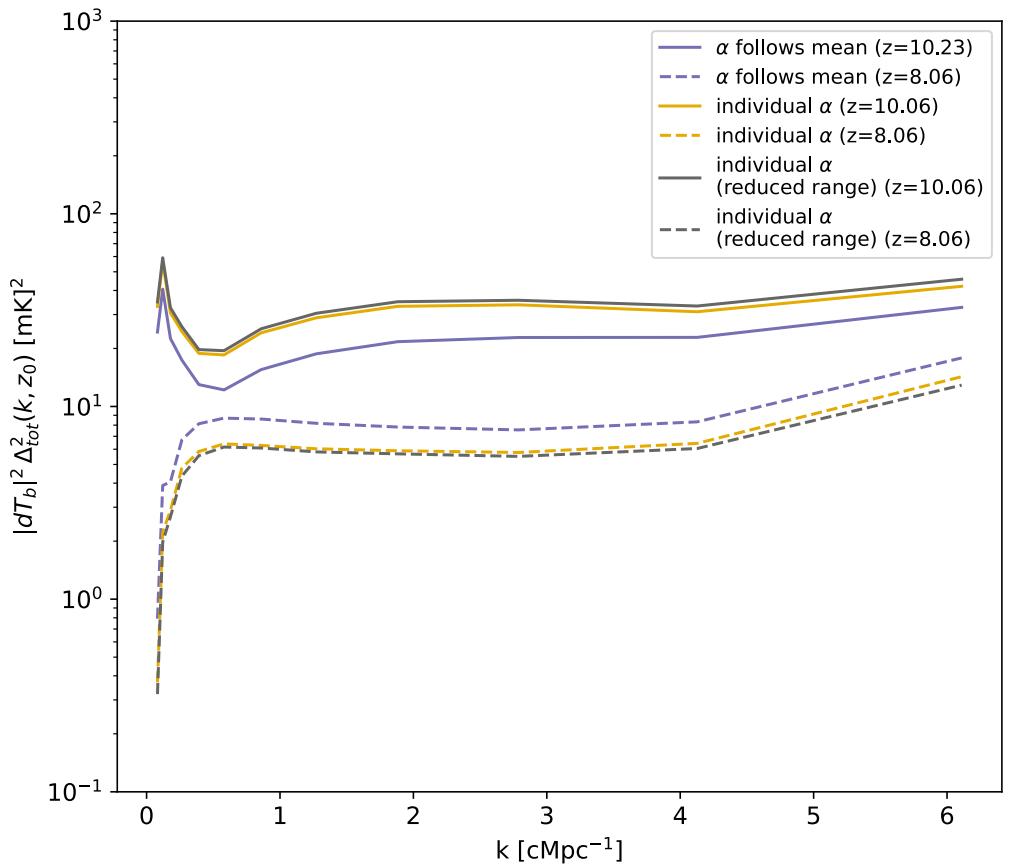
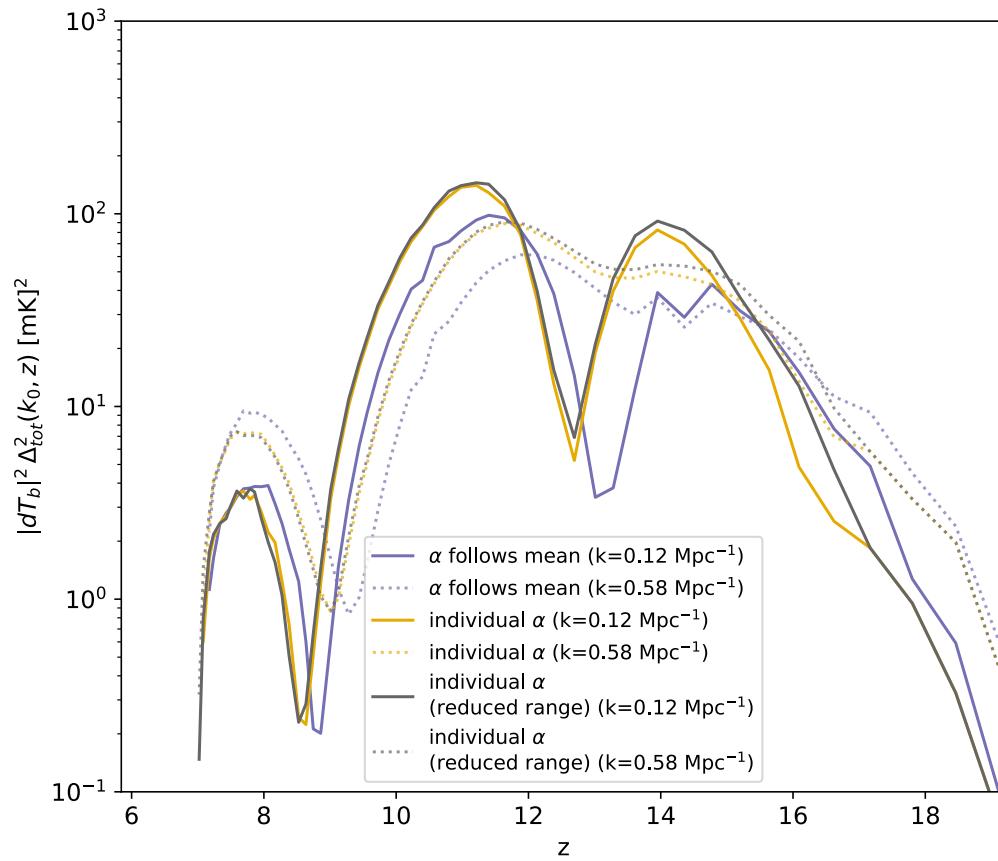
Map outputs

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- richer structures due to combined effects
- clear distinction between “foreground” and “background” effects





Simulating the
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Conclusion

- Summary
- Outlook

- BEoRN a semi-numerical tool to simulate the 21-cm signal
 - uses the *halo model of reionization*
 - describes sources in terms of their host DM halo
 - ⇒ central dependence on halo growth
- more accurate treatment of **individual** mass accretion
 - leads to significant changes to reionization history
 - map-level fluctuations
- BEoRN python package: <https://github.com/cosmic-reionization/beorn>
 - simulation-agnostic
 - easier to use
 - fully parallelized

- further validation
- investigation + parameterization of stochasticity
- application to larger volumes

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End

— Thank you for your attention

References

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- V. Springel *et al.*, “Simulations of the formation, evolution and clustering of galaxies and quasars,” *Nature*, vol. 435, no. 7042, pp. 629–636, Jun. 2005, doi: 10.1038/nature03597.

$$\rho_\alpha(r \mid M, z) = \frac{(1+z)^2}{4\pi r^2} \cdot \sum_{n=2}^{n_m} f_n \cdot \varepsilon_\alpha(\nu') \cdot f_\star \cdot \dot{M}(z' \mid M, z)$$

with the lookback redshift z' so $\nu' = \nu \cdot (1+z')/(1+z)$

⇒ coupling coefficient

$$x_\alpha(r \mid M, z) = \frac{1.81 \cdot 10^{11}}{1+z} \cdot S_\alpha(z) \cdot \rho_\alpha(r \mid M, z)$$

with a suppression factor $S_\alpha(z)$

$$\rho_{\text{xray}}(r \mid M, z) = \frac{1}{r^2} \sum_i f_i f_{X,h} \cdot \int_{\nu_{\text{th}}^i}^{\infty} d\nu (\nu - \nu_{\text{th}}^i) h_P \sigma_i(\nu) e^{-\tau_\nu} f_\star \dot{M}(z' \mid M, z)$$

$$\Rightarrow \frac{3}{2} \cdot \frac{d\rho_h(r \mid M, z)}{dz} = \frac{3\rho_h(r \mid M, z)}{1+z} - \frac{\rho_{\text{xray}}(r \mid M, z)}{k_B(1+z)H(z)}$$

with the Boltzmann constant k_B and $H(z)$ is the Hubble parameter

The comoving ionized volume around a source of ionizing photons satisfies the differential equation

$$\frac{dV}{dt} = \frac{\dot{N}_{\text{ion}(t)}}{\bar{n}_H^0} - \alpha_B \cdot \frac{C}{a^3} \cdot \bar{n}_H^0 \cdot V$$

bubble radius $R_b = \sqrt[3]{\frac{3}{4\pi}V(M, z)}$ and using the Heaviside step function θ_H :

$$x_{\text{HII}}(r \mid M, z) = \theta_H[R_b(M, z) - r]$$

Validation

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